

Determinate Values for Quantum Observables

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Abstract

This is a comment on J. A. Barrett’s article “The Preferred-Basis Problem and the Quantum Mechanics of Everything” in *Brit. J. Phil. Sci.* **56** (2005), which concerns theories postulating that certain quantum observables have determinate values, corresponding to additional (often called “hidden”) variables. I point out that it is far from clear, for most observables, what such a postulate is supposed to mean, unless the postulated additional variable is related to a clear ontology in space-time, such as particle world lines, string world sheets, or fields.

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In his recent article (Barrett 2005), Jeffrey A. Barrett developed an astute analysis of the problems that would arise for Bohmian mechanics if mental states did not supervene on the positions of the particles constituting the brain. My comment on his article is not so much a criticism but rather concerns a point that I think should be kept in mind in this context but that Barrett did not mention in his article. The point is that one is not always free to add a postulate like “there exist determinate values for the quantum-mechanical observable Λ ” because it is far from clear, for most observables, what such a postulate would *mean*. It is the purpose of this comment to elucidate this point. Related considerations can be found in (Daumer et al. 1996), in particular Section 6, and in (Allori et al. 2006), in particular Sections 4–8.

According to Bohmian mechanics (also known as de Broglie–Bohm pilot-wave theory) (Bohm 1952; Bell 1966; Goldstein 2001) there are point particles moving in 3-space, which have a determinate position at every time, in such a way that the configuration q is random with probability distribution $|\psi_t(q)|^2$ at every time t . In the language of quantum mechanics, one could say that the position observables have determinate values in Bohmian mechanics, namely the actual positions of the particles, while other observables are not attributed determinate values (except in special cases, such as that of the wave function being an eigenfunction of the observable). Similarly, Bohm (1952) in his quantum field theory model postulated determinate values of the electromagnetic field operators, and Bell (1986) in his quantum field theory model postulated determinate values of the fermion number operators. Since one can understand in all of these cases what it means for these observables to have determinate values, one might get the

impression that one can always postulate the existence of such additional (“hidden”) variables that provide the actual values for any observable Λ one may wish to choose. However, that means to overlook the following problem.

While it is clear what it means to postulate that there are point particles moving in 3-space, it is not clear at all what it should mean to postulate that there is a variable λ , governed by certain equations, until we specify the relation of λ to the physical ontology in space-time, such as particle world lines or string world sheets or fields. The relation could be that λ is a function of the ontology in space-time, or that λ influences its behavior; e.g., the velocity is a function of the particle world line, and the mass affects the world line as a parameter in the equation of motion (in Bohmian or classical mechanics).

Let me approach the problem by a sequence of examples. Suppose first that some theory, some alternative to Bohmian mechanics, postulates that not the positions, but the velocities are the observables that have determinate values.¹ Then this postulate could be understood as saying that

$$\begin{aligned} &\text{there are point particles moving in 3-space whose velocities} \\ &v_i \text{ have joint probability distribution } |\tilde{\psi}_t(\dots m_i v_i \dots)|^2 \end{aligned} \tag{1}$$

where $\tilde{\psi}$ is the momentum representation (Fourier transform) of ψ and m_i the mass of particle i . The point is that I did not simply say,

$$\begin{aligned} &\text{there are variables } v_i \text{ with joint probability} \\ &\text{distribution } |\tilde{\psi}_t(\dots m_i v_i \dots)|^2. \end{aligned} \tag{2}$$

I would not be able to understand the meaning of (2) without further information, whereas (1) is clear. As an example of the ambiguity of (2), I could as well postulate

$$\begin{aligned} &\text{there are point particles moving in 3-space whose **positions**} \\ &v_i \text{ have joint probability distribution } |\tilde{\psi}_t(\dots m_i v_i \dots)|^2. \end{aligned} \tag{3}$$

This would seem compatible with (2), indeed as compatible as (1).² This fact indicates that (2) is an incomplete description of a theory and needs to be complemented by either (1), or (3), or a similar statement.

Suppose now, as a second example, that some theory postulates that

$$\text{“energy is the observable with determinate values.”} \tag{4}$$

¹The reader might think at this point that such a theory would not solve the measurement problem; I agree, but this is not the point I want to make now. The reader might also think that the way Bohmian mechanics ensures that the probability distribution is $|\psi_t(q)|^2$ at all times t , cannot be transferred to the momentum representation; again, I agree, but this is not the point I want to make now.

²The reader might think that such a theory would not make the same predictions as quantum mechanics; I agree, but this is not the point I want to make now.

In analogy to the previous example, in particular to (1), this could be understood as saying that

$$\begin{aligned} &\text{there are point particles moving in 3-space whose} \\ &\text{positions } q_i \text{ and velocities } v_i \text{ have such a distribution } \mathbb{P} \\ &\text{that } \mathbb{P}\left(\sum_i \frac{1}{2}m_i v_i^2 + V(\dots q_i \dots) = E\right) = \langle \psi | P_E | \psi \rangle \end{aligned} \quad (5)$$

with V the potential and P_E the projection to the eigenspace of the Hamiltonian with eigenvalue E . But it could also be understood in a different way, for example as saying that

$$\begin{aligned} &\text{there is a field } \phi(x, t) \text{ on space-time whose distribution } \mathbb{P} \\ &\text{is such that } \mathbb{P}\left(\int d^3x |\nabla \phi(x, t)|^2 = E\right) = \langle \psi | P_E | \psi \rangle. \end{aligned} \quad (6)$$

Moreover, (5) or (6) is not what comes to mind when one reads (4). What comes to mind, instead, is that there are no point particles (and no strings and no fields), but there is only *energy as such*. And this is the step where it is not clear, at least to me, what the statement is supposed to mean.

The difficulty is connected to the question, What is the meaning of the word “energy”? In classical mechanics or field theory, energy is defined to be $\sum_i \frac{1}{2}m_i v_i^2 + V(\dots q_i \dots)$ respectively $\int d^3x |\nabla \phi(x, t)|^2$. If I am supposed to understand the word “energy” without any further definition, I am lost. For example, I wonder what would then be the difference if the same determinate value λ were not *energy as such*, but, say, *charge as such*? Moreover, part of the physicist’s psychological expectation connected with the statement that “energy has the value λ ” is that a measurement of the energy will have result λ . But *that* cannot be postulated, but either follows or does not follow from the macroscopic motion of matter constituting the measurement apparatus; and this matter should correspond to some ontology in space-time.

As a third example, let us consider various ways of dealing with spin in theories like Bohmian mechanics. The modern way of treating spin, which goes back to Bell (1966), postulates no further additional variable for spin in addition to the particle positions; spin is taken into account by assuming that the wave function of an electron is \mathbb{C}^2 -valued, $\psi(q) = (\psi_\uparrow(q), \psi_\downarrow(q))$, and the probability distribution of its position is $\|\psi(q)\|^2 = |\psi_\uparrow(q)|^2 + |\psi_\downarrow(q)|^2$ with $\|\cdot\|$ the norm in \mathbb{C}^2 . Another possible approach that has been investigated by Bohm and Hiley (1993) and Holland (1995) assumes that an electron is a tiny rotating sphere, and that the spin observables are somehow connected to the actual speed of rotation. While there are reasons to prefer Bell’s approach, the relevant aspect here is that it is understandable what “speed of rotation” means if an electron is an extended body with an orientation in space.³ Let me contrast this with a

³A side remark concerning a subtlety: There is an ambiguity about what it means to be a sphere. One possibility, providing what is needed to make sense of the notion of “speed of rotation,” is that every point on the sphere has a trajectory, identifying points on the sphere at different times. The other possibility is to postulate that the sphere is simply a subset of space filled with matter, without any identification of points at different times. In the latter view, the speed of rotation is not defined.

postulate that is not understandable, exactly for the reason that it introduces a variable λ without relating it to the constituents of the physical ontology. The fact that the wave function of an electron can be written as a function $\psi(q, s)$ of a continuous (position) variable $q \in \mathbb{R}^3$ and a discrete (spin) variable $s \in \{\uparrow, \downarrow\}$ may suggest to introduce, in addition to the particle position, another additional variable $\lambda \in \{\uparrow, \downarrow\}$ representing the “actual (z -component of the) spin.”⁴ My objection is that *it does not make sense* to merely postulate “there is a variable λ which can assume the values \uparrow and \downarrow and obeys these-and-these laws.” To make it meaningful, one might postulate that an electron is not a point particle but, say, a tiny (asymmetric) rod that always points in either the z direction or the $-z$ direction, and that λ indicates this direction; or one might specify how λ influences the particle’s trajectory in space.

Many readers may have the following question at this point: Would it not be sufficient for a newly introduced variable λ if there was a psycho-physical law according to which my conscious experience is a certain function of λ ? Would it not be unnecessary then to relate λ to some physical ontology in space and time? To be sure, since such a theory would predict the same experience, it is irrefutable; but it has the flavor of a brain-in-the-vat scenario. According to such a theory, the physical world would not at all be like what we usually think it is, and that is why I do not seriously consider this possibility. The world would not contain any trees, or rocks, or chairs, as it would not be able to contain different amounts of matter in different regions of space and time. Instead, there would just be the abstract quantity λ , giving rise to illusory experiences of trees and rocks and chairs. I cannot stop believing that there are trees and rocks and chairs.

I hope that these examples help convey why I think one should regard the postulate “the observable Λ has determinate value λ ” as incomprehensible in a physical theory unless λ is connected to a physical ontology in space-time such as particle world lines or string world sheets or fields. Such a connection is naturally suggested for some observables together with certain ontologies, such as position operators with particle ontology or field operators with field ontology. For the velocity operator $-\frac{i}{m}\nabla$, a connection is less strongly suggested, and for the energy operator H hardly at all. For most operators no connection suggests itself (a concrete example is $\Lambda = q - i\nabla$).

For Barrett’s (2005) discussion this means that “ E -theory” (according to which all quantum-mechanical observables have determinate values) appears not to be among the viable options, i.e., among the meaningful theories, as for most observables Λ the postulate that Λ has determinate values λ remains incomprehensible. Whether “ Q -theory” (according to which the observable Q on which mental states supervene, whichever it may turn out to be, has determinate values) is a viable option seems to depend on which observable Q turns out to be. For Bohmian mechanics, our discussion means that it is misleading to say, “the position observables have determinate values.” A more appropriate description of Bohmian mechanics is, “there are particles moving in 3-space whose positions have the distribution of the quantum-mechanical position observables.”

⁴The reader might think at this point that such a postulate would unduly prefer the z -axis over other directions; I agree, but this is not the point I want to make now.

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References

- Allori, V., Goldstein, S., Tumulka, R. and Zanghì, N. [2006]: ‘On the Common Structure of Bohmian Mechanics and the Ghirardi–Rimini–Weber Theory’, preprint <http://arxiv.org/quant-ph/0603027>
- Barrett, J. A. [2005]: ‘The Preferred-Basis Problem and the Quantum Mechanics of Everything’, *British Journal for the Philosophy of Science*, **56**, pp. 199–220.
- Bell, J. S. [1966]: ‘On the problem of hidden variables in quantum mechanics’, *Reviews of Modern Physics*, **38**, pp. 447–52. Reprinted in (Bell 1987).
- Bell, J. S. [1986]: ‘Beables for quantum field theory’, *Physics Reports*, **137**, pp. 49–54. Reprinted in (Bell 1987).
- Bell, J. S. [1987]: *Speakable and unspeakable in quantum mechanics*, Cambridge: Cambridge University Press.
- Bohm, D. [1952]: ‘A Suggested Interpretation of the Quantum Theory in Terms of “Hidden” Variables, I and II’, *Physical Review*, **85**, pp. 166–93.
- Bohm, D. and Hiley, B. J. [1993]: *The Undivided Universe: An Ontological Interpretation of Quantum Theory*, London: Routledge.
- Daumer, M., Dürr, D., Goldstein, S. and Zanghì, N. [1996]: ‘Naive Realism about Operators’, *Erkenntnis*, **45**, pp. 379–97. Preprint <http://arxiv.org/quant-ph/9601013>
- Goldstein, S. [2001]: ‘Bohmian mechanics’, in E. N. Zalta (ed.), *Stanford Encyclopedia of Philosophy*, published online by Stanford University at <http://plato.stanford.edu/entries/qm-bohm/>
- Holland, P. R. [1995]: *The quantum theory of motion: An account of the de Broglie-Bohm causal interpretation of quantum mechanics*, Cambridge: Cambridge University Press.